

QCD matter in extreme conditions

Gergely Endrődi

University of Bielefeld



**UNIVERSITÄT
BIELEFELD**

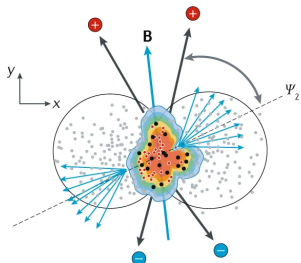
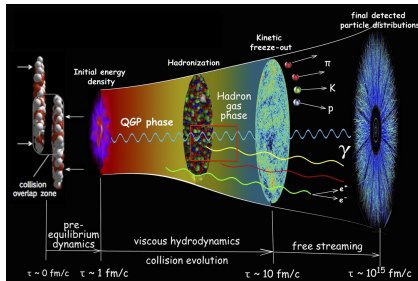


KHuK Annual Meeting
Bad Honnef, December 9, 2022

Introduction

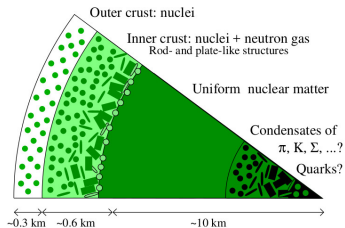
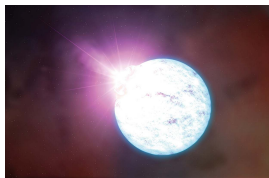
Quarks and gluons in extreme conditions

- ▶ heavy ion collisions $T \lesssim 10^{12} \text{ }^\circ\text{C} = 200 \text{ MeV}$, $n \lesssim 0.12 \text{ fm}^{-3}$
 $B \lesssim 10^{19} \text{ G} = 0.3 \text{ GeV}^2/e$



Quarks and gluons in extreme conditions

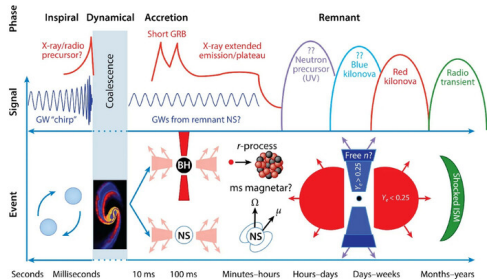
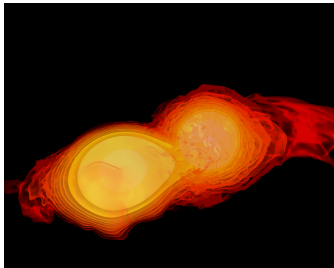
- ▶ heavy ion collisions $T \lesssim 10^{12} \text{ }^\circ\text{C} = 200 \text{ MeV}$, $n \lesssim 0.12 \text{ fm}^{-3}$
 $B \lesssim 10^{19} \text{ G} = 0.3 \text{ GeV}^2/e$
- ▶ neutron stars $T \lesssim 1 \text{ keV}$, $n \lesssim 2 \text{ fm}^{-3}$
magnetars $B \lesssim 10^{15} \text{ G}$



 Lattimer, Nature Astronomy 2019

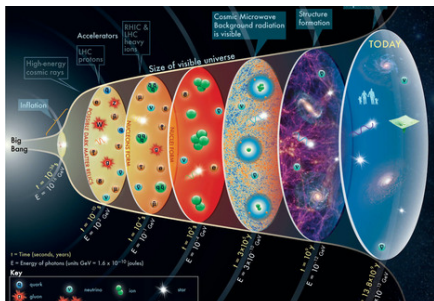
Quarks and gluons in extreme conditions

- ▶ heavy ion collisions $T \lesssim 10^{12} \text{ }^\circ\text{C} = 200 \text{ MeV}$, $n \lesssim 0.12 \text{ fm}^{-3}$
 $B \lesssim 10^{19} \text{ G} = 0.3 \text{ GeV}^2/e$
- ▶ neutron stars $T \lesssim 1 \text{ keV}$, $n \lesssim 2 \text{ fm}^{-3}$
magnetars $B \lesssim 10^{15} \text{ G}$
- ▶ neutron star mergers $T \lesssim 50 \text{ MeV}$

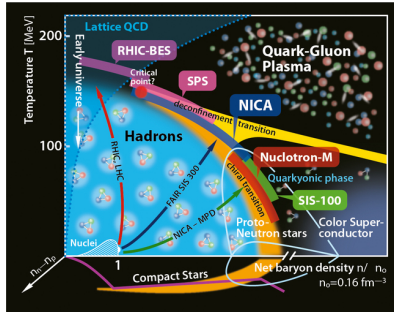
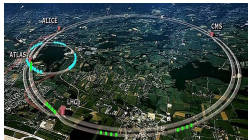


Quarks and gluons in extreme conditions

- ▶ heavy ion collisions $T \lesssim 10^{12} \text{ }^\circ\text{C} = 200 \text{ MeV}$, $n \lesssim 0.12 \text{ fm}^{-3}$
 $B \lesssim 10^{19} \text{ G} = 0.3 \text{ GeV}^2/e$
- ▶ neutron stars $T \lesssim 1 \text{ keV}$, $n \lesssim 2 \text{ fm}^{-3}$
magnetars $B \lesssim 10^{15} \text{ G}$
- ▶ neutron star mergers $T \lesssim 50 \text{ MeV}$
- ▶ early universe, QCD epoch $T \lesssim 200 \text{ MeV}$
standard scenario: $n \approx 0$



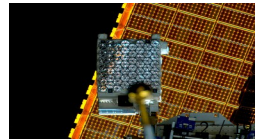
Major experimental and observational campaigns



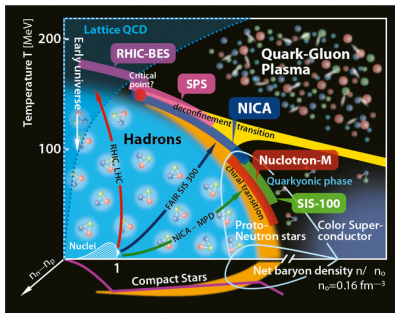
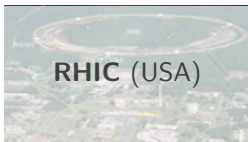
Observational astronomy



Heavy ion collisions



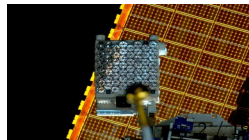
Major experimental and observational campaigns



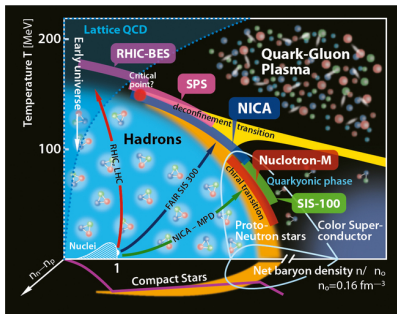
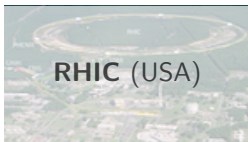
Observational astronomy



Heavy ion collisions



Major experimental and observational campaigns



Observational astronomy



Heavy ion collisions



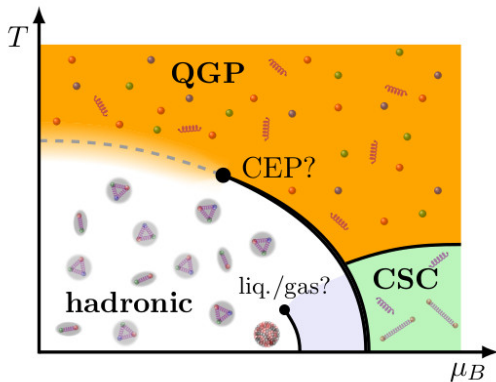
QCD phase diagram(s)

Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$

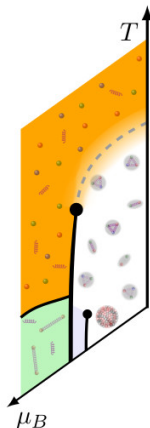
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram



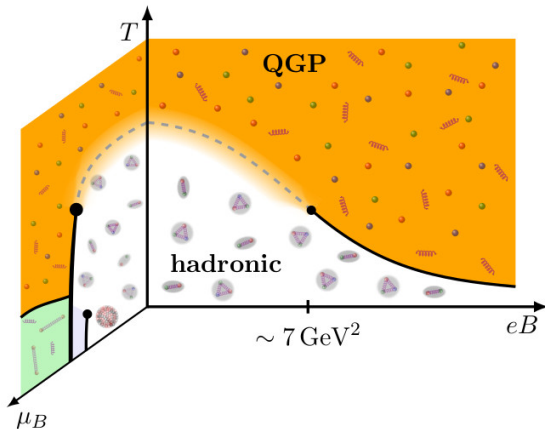
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



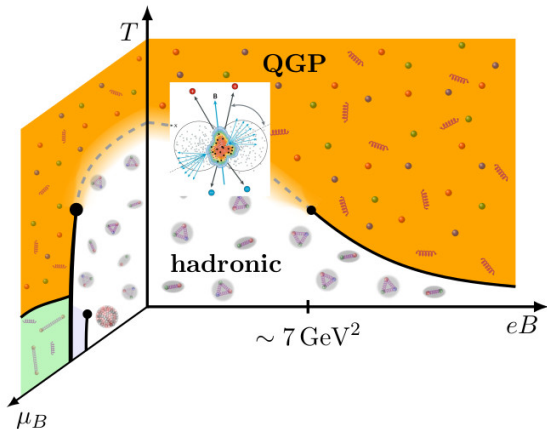
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



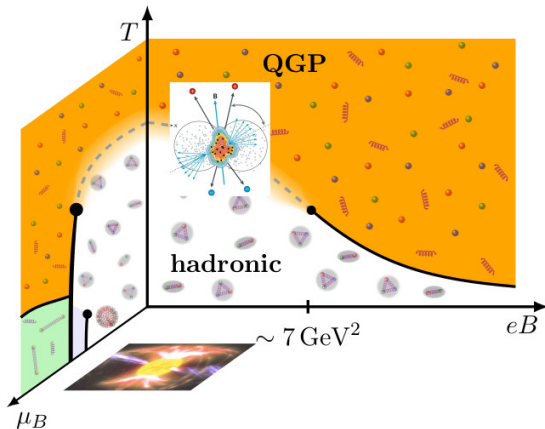
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



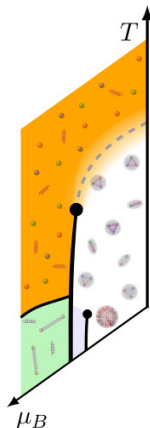
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



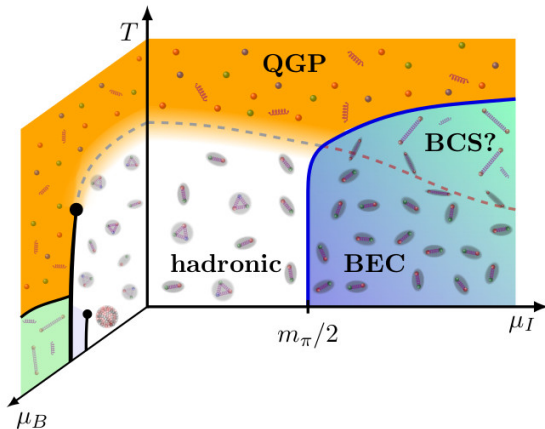
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



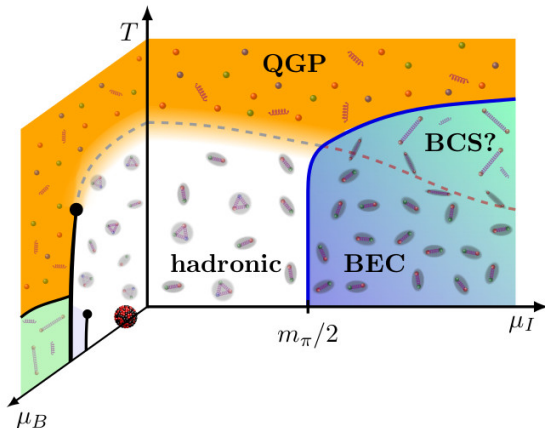
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



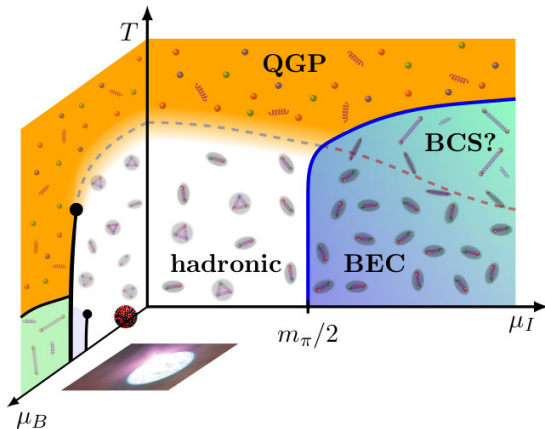
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



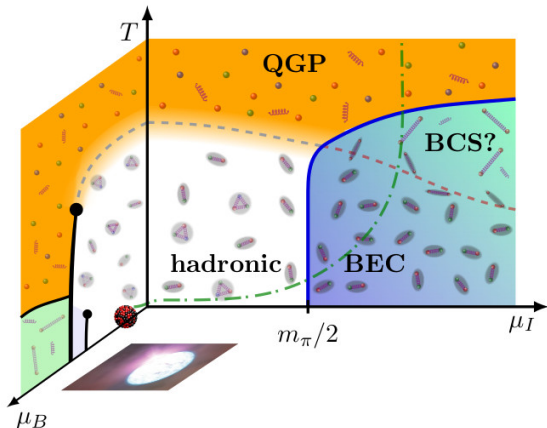
Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



Phase diagram

- ▶ control parameters: $T, n \leftrightarrow \mu, B$
 $\mu_{\{u,d,s\}} / \mu_{\{B,Q,S\}} / \mu_{\{B,I,S\}}$
- ▶ well-known famous phase diagram
- ▶ well-known, less famous phase diagrams



Methods to study QCD thermodynamics

Lattice simulations

- ▶ Euclidean QCD path integral over gauge field \mathcal{A}

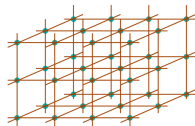
$$\mathcal{Z} = \int \mathcal{D}\mathcal{A} e^{-S_g[\mathcal{A}]} \det[\not{D}[\mathcal{A}] + m]$$

- ▶ Monte-Carlo simulations need: $\det[\not{D} + m] \in \mathbb{R}^+$

need Γ : $\Gamma \not{D} \Gamma^\dagger = \not{D}^\dagger, \quad \Gamma^\dagger \Gamma = 1$

- ▶ $\exists \Gamma$ ✓

$$B, \mu_I, i\mu_B, iE$$

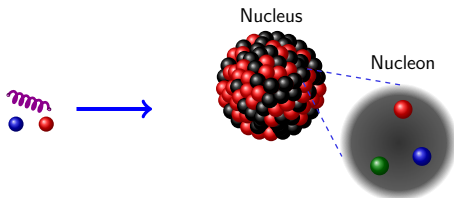


continuum limit to recover full theory

- ▶ $\nexists \Gamma$ \times complex action (sign) problem
 μ_B, E

Functional renormalization group

- ▶ renormalization group flow from UV to IR
✍ Kadanoff '66 ✍ Wilson '71
- ▶ for QCD: from quarks and gluons to hadrons and nuclei



via successive integration of high-momentum modes
Wetterich equation ✍ Wetterich '92

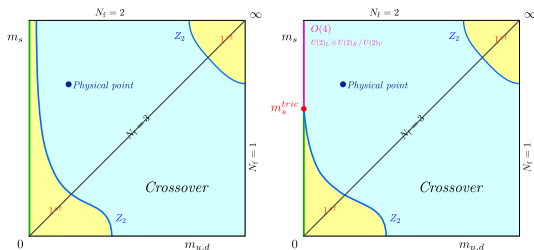
- ▶ exact flow equation, access to complete phase diagram ✓
- ▶ requires approximations (truncations, Ansätze) to solve ✗

Thermodynamics at $\mu_B = 0$

Chiral limit at zero density

- ▶ transition at physical quark masses is a crossover

✍ Aoki et al. '06 ✍ Bhattacharya et al. '14

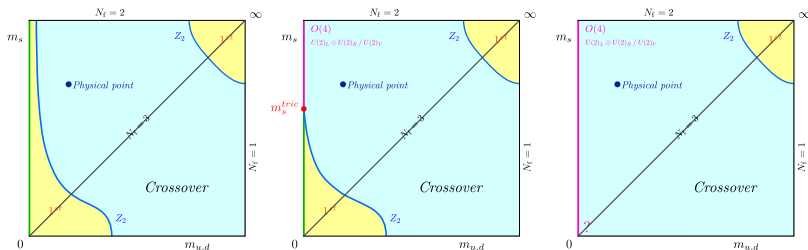


- ▶ chiral limit: expect $1_{N_f=3}^{st}$ and $1_{N_f=2}^{st} / 2_{N_f=2}^{nd}$ depending on $U_A(1)$ restoration ✍ Pisarski, Wilczek '84

Chiral limit at zero density

- ▶ transition at physical quark masses is a crossover

🔗 Aoki et al. '06 🔗 Bhattacharya et al. '14



- ▶ chiral limit: expect $1^{st}_{N_f=3}$ and $1^{st}_{N_f=2}/2^{nd}_{N_f=2}$ depending on $U_A(1)$ restoration 🔗 Pisarski, Wilczek '84

- ▶ lattice exploiting tricritical scaling in N_f : $2^{nd}_{N_f=2,3}$

🔗 Cuteri, Philipsen, Sciarra '21

FRG including 't Hooft coupling $2^{nd}_{N_f=2}$ 🔗 Braun et al. '20

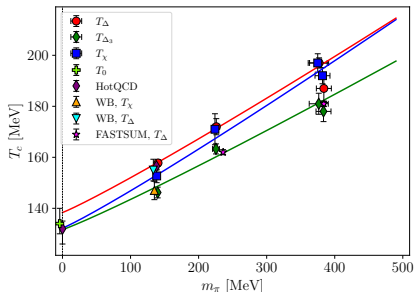
Chiral limit at zero density

- scaling of pseudocritical temperature gives:

$$T_c(m_{ud} = 0, m_s^{\text{phys}}) = 132_{-6}^{+3} \text{ MeV} \quad \text{Ding et al. '19}$$

compare [Kotov et al. '21](#) [Borsányi et al. '20](#) [Aarts et al. '20](#)

$$N_f = 3: T_c(m_{ud} = m_s = 0) = 98_{-6}^{+3} \text{ MeV} \quad \text{Dini et al. '21}$$



Chiral limit at zero density

- scaling of pseudocritical temperature gives:

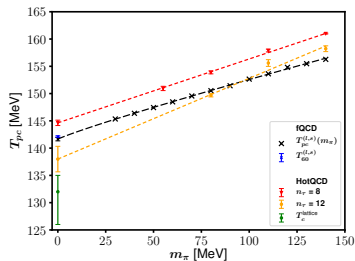
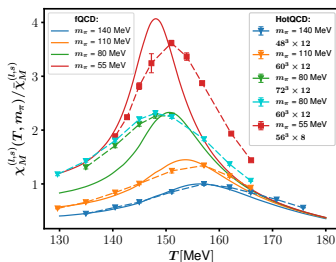
$$T_c(m_{ud} = 0, m_s^{\text{phys}}) = 132_{-6}^{+3} \text{ MeV} \quad \text{Ding et al. '19}$$

compare [Kotov et al. '21](#) [Borsányi et al. '20](#) [Aarts et al. '20](#)

$$N_f = 3: T_c(m_{ud} = m_s = 0) = 98_{-6}^{+3} \text{ MeV} \quad \text{Dini et al. '21}$$

- direct comparison between FRG and lattice

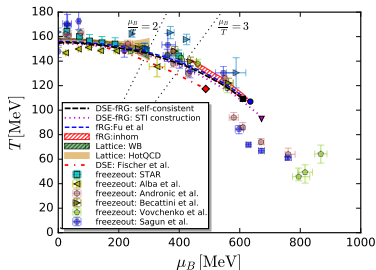
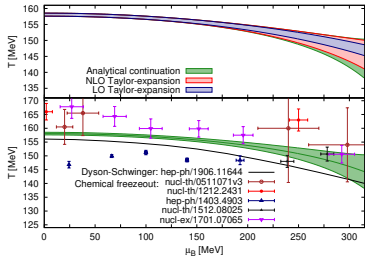
[Braun et al. '20](#) [Ding et al. '19](#)



Thermodynamics at $\mu_B > 0$

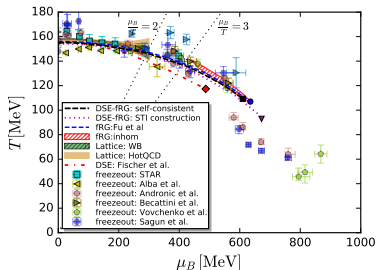
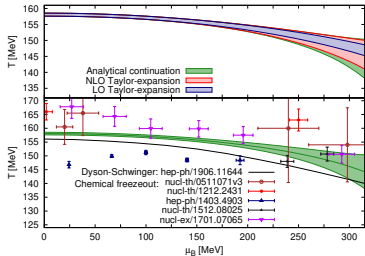
Phase diagram in the $T - \mu_B$ plane

- ▶ analytical continuation of lattice results at $i\mu_B > 0$
consistency with Taylor expansion [Borsányi et al. '20](#)
- ▶ functional methods prefer critical endpoint
- ▶ FRG: [Fu, Pawłowski, Rennecke '20](#) [Gao, Pawłowski '20](#)
[Otto, Busch, Schaefer '22](#)
- ▶ DSE: including meson backcoupling effects [Gunkel, Fischer '21](#)



Phase diagram in the $T - \mu_B$ plane

- ▶ analytical continuation of lattice results at $i\mu_B > 0$
consistency with Taylor expansion [Borsányi et al. '20](#)
- ▶ functional methods prefer critical endpoint
- ▶ FRG: [Fu, Pawłowski, Rennecke '20](#) [Gao, Pawłowski '20](#)
[Otto, Busch, Schaefer '22](#)
- ▶ DSE: including meson backcoupling effects [Gunkel, Fischer '21](#)



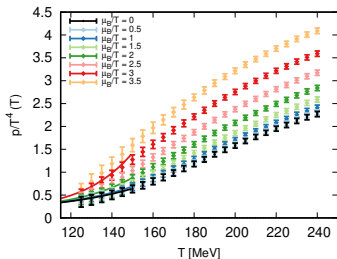
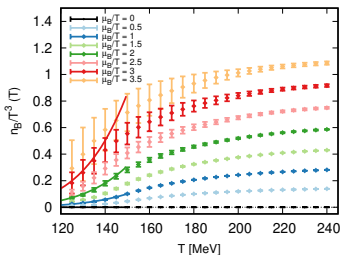
- ▶ inhomogeneous instability at large μ_B ?

Equation of state

- ▶ combining Taylor expansion in μ_B and shift in T

$$\mathcal{O}(T, \mu_B) \approx \mathcal{O}(T - \kappa\mu_B^2, 0)$$

- ▶ primary observable: baryon density



 Borsányi et al. '21

- ▶ at zero strangeness density, relevant for HIC  Borsányi et al. '22

Further results

- ▶ alternative resummation schemes [✍ Mondal et al. '21](#)
- ▶ imaginary chemical potentials and Roberge-Weiss phase transitions [✍ Brandt et al. '22](#)
- ▶ QCD transition in the heavy quark/quenched limit
[✍ Borsányi et al. '22](#)
- ▶ thermal effects on hadrons, chiral-spin symmetry
[✍ Aarts et al. '20](#) [✍ Glozman, Philipsen, Pisarski '22](#)
- ▶ transport properties - photon emissivity [✍ Cé et al. '22](#)
- ▶ heavy quark diffusion [✍ Brambilla et al. '22](#) [✍ Altenkort et al. '22](#)

Thermodynamics at $\mu_I > 0$

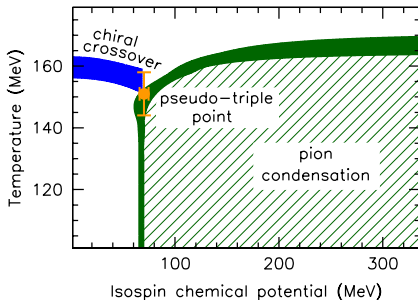
Phase diagram in the $T - \mu_I$ plane

▶ $\mu_I = \mu_u - \mu_d$

▶ phases:

hadronic, quark-gluon plasma, BEC of charged pions

✎ Brandt, Endrődi, Schmalzbauer '17 ✎ Brandt, Endrődi '19



▶ compares well to χ PT at low T

✎ Son, Stephanov '01 ✎ Adhikari et al. '20

Equation of state on the lattice

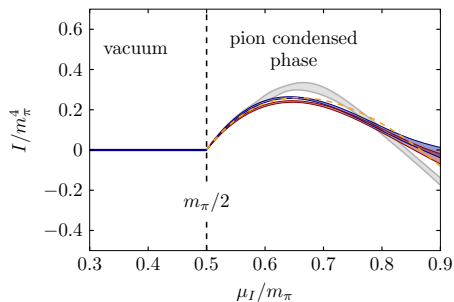
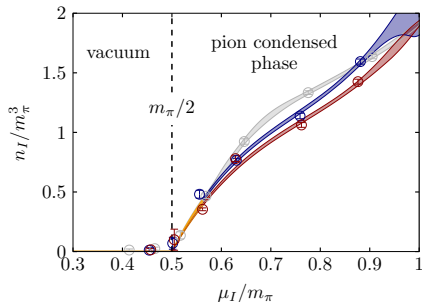
- ▶ primary observable: isospin density

$$n_I = \frac{T}{V} \frac{\partial \log \mathcal{Z}}{\partial \mu_I}, \quad p(T, \mu_I) - p(T, 0) = \int_0^{\mu_I} d\mu'_I n_I(\mu'_I)$$

- ▶ results at $T \approx 0$

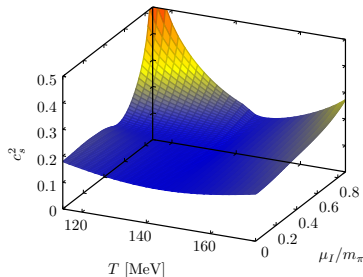
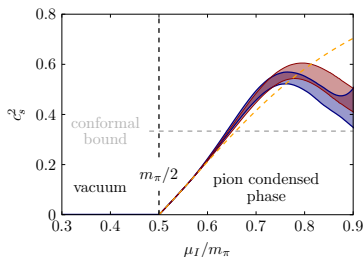
✍ Brandt, Endrődi, Fraga, Hippert, Schaffner-Bielich, Schmalzbauer '18

✍ Brandt, Cuteri, Endrődi '22



Equation of state on the lattice

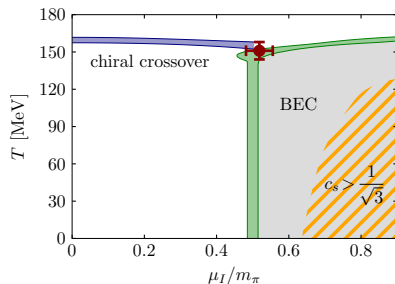
- ▶ results at $T \neq 0$ ✍ Brandt, Cuteri, Endrődi '22
✍ Vovchenko, Brandt, Cuteri, Endrődi, Hajkarim, Schaffner-Bielich '20
- ▶ interaction measure peak shifts to lower T as μ_I grows
- ▶ speed of sound **above $1/\sqrt{3}$** at high μ_I and intermediate T



- ▶ EoS gets very stiff inside pion condensation phase

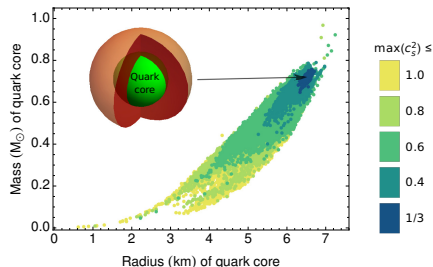
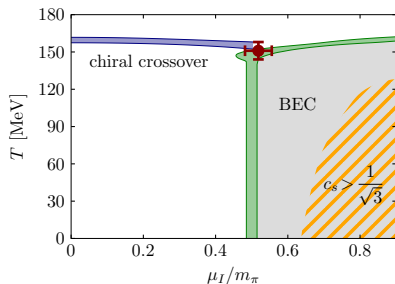
Speed of sound

- ▶ 'supersonic' region of pion condensate
- ▶ first time that $c_s > 1/\sqrt{3}$ found in a first-principles lattice QCD calculation



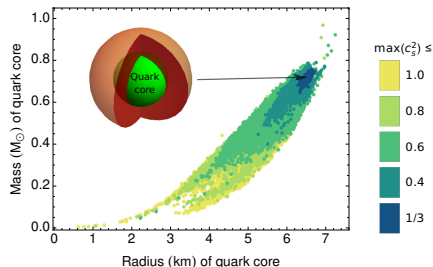
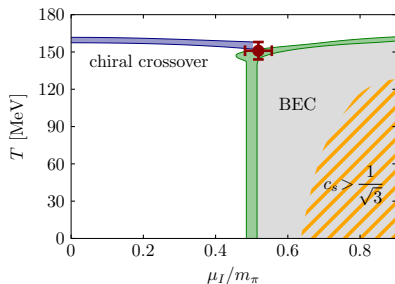
Speed of sound

- ▶ 'supersonic' region of pion condensate
- ▶ first time that $c_s > 1/\sqrt{3}$ found in a first-principles lattice QCD calculation
- ▶ relevance of c_s for neutron star modeling [Annala et al. '19](#)



Speed of sound

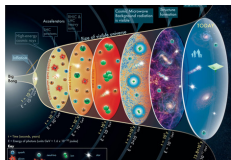
- ▶ 'supersonic' region of pion condensate
- ▶ first time that $c_s > 1/\sqrt{3}$ found in a first-principles lattice QCD calculation
- ▶ relevance of c_s for neutron star modeling [Annala et al. '19](#)
- ▶ c_s at $\mu_B > 0$ from FRG and χ EFT
[Braun, Schallmo '22](#) [Leonhardt et al. '20](#)



Cosmological implications

Cosmic trajectories

- ▶ early Universe



- ▶ conservation equations for isentropic expansion

$$\frac{n_B}{s} = b, \quad \frac{n_Q}{s} = 0, \quad \frac{n_{L\alpha}}{s} = l_\alpha \quad (\alpha \in \{e, \mu, \tau\})$$

- ▶ parameters: T , μ_B , μ_Q , $\mu_{L\alpha}$
- ▶ experimental constraints [Planck coll. '15](#) [Oldengott, Schwarz '17](#)

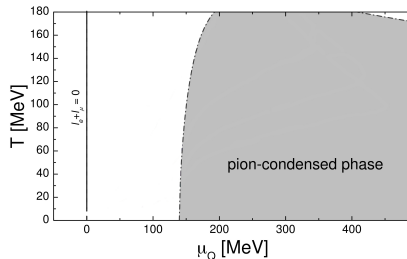
$$b = (8.60 \pm 0.06) \cdot 10^{-11}, \quad |l_e + l_\mu + l_\tau| < 0.012$$

(the individual l_α may have opposite signs)

- ▶ $n_Q = 0$ with $l_e > 0$ allows equilibrium of e^- , ν_e , π^+

Cosmic trajectories

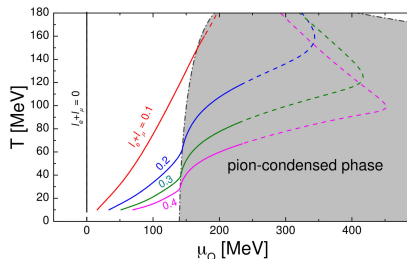
- ▶ cosmic trajectory $T(\mu_Q)$ is solved for
- ▶ standard scenario ($I_\alpha = 0$): $\mu_Q = 0$ for all T



✍ Vovchenko, Brandt, Cuteri, Endrődi, Hajkarim, Schaffner-Bielich '20

Cosmic trajectories

- ▶ cosmic trajectory $T(\mu_Q)$ is solved for
- ▶ standard scenario ($I_\alpha = 0$): $\mu_Q = 0$ for all T



✍ Vovchenko, Brandt, Cuteri, Endrődi, Hajkarim, Schaffner-Bielich '20

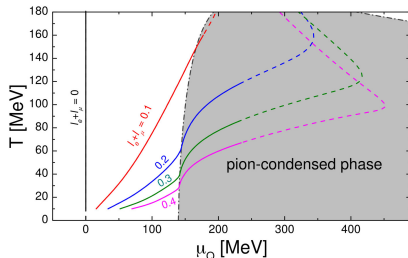
- ▶ condition for pion condensation to occur:

$$|I_e + I_\mu + I_\tau| < 0.012$$

$$|I_e + I_\mu| \gtrsim 0.1$$

Cosmic trajectories

- ▶ cosmic trajectory $T(\mu_Q)$ is solved for
- ▶ standard scenario ($I_\alpha = 0$): $\mu_Q = 0$ for all T



✍ Vovchenko, Brandt, Cuteri, Endrődi, Hajkarim, Schaffner-Bielich '20

- ▶ condition for pion condensation to occur:

$$|I_e + I_\mu + I_\tau| < 0.012$$

$$|I_e + I_\mu| \gtrsim 0.1$$

- ▶ enhanced primordial grav. waves (SKA)

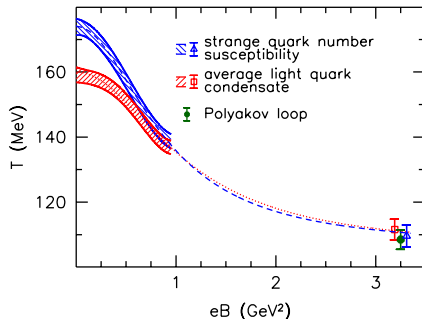


Thermodynamics at $B > 0$

Phase diagram and critical point

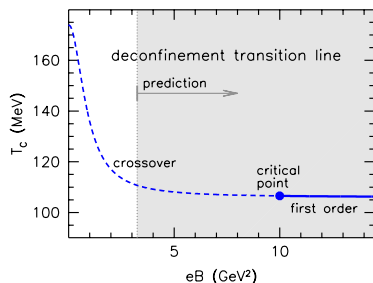
- ▶ physical m_π , staggered quarks, continuum limit

✍ Bali, Bruckmann, Endrődi, Fodor, Katz et al. '11 ✍ '12 ✍ Endrődi '15



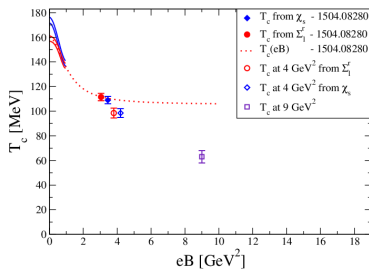
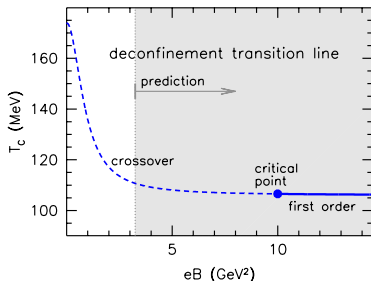
Phase diagram and critical point

- ▶ physical m_π , staggered quarks, continuum limit
✍ Bali, Bruckmann, Endrődi, Fodor, Katz et al. '11 ✍ '12 ✍ Endrődi '15
- ▶ transition strengthens \Rightarrow critical point at $eB_c \approx 10(2) \text{ GeV}^2$
✍ Endrődi '15



Phase diagram and critical point

- ▶ physical m_π , staggered quarks, continuum limit
✍ Bali, Bruckmann, Endrödi, Fodor, Katz et al. '11 ✍ '12 ✍ Endrödi '15
- ▶ transition strengthens \Rightarrow critical point at $eB_c \approx 10(2) \text{ GeV}^2$
✍ Endrödi '15
- ▶ simulating up to $eB \approx 9 \text{ GeV}^2 \Rightarrow 4 \text{ GeV}^2 < eB_c < 9 \text{ GeV}^2$
✍ D'Elia, Maio, Sanfilippo, Stanzione '21



Further results on magnetic fields

- ▶ fluctuations of conserved charges at $B > 0$, $T > 0$
✍ Ding et al. '21
- ▶ anomalous transport phenomena at $B > 0$
✍ Astrakhantsev et al. '20 ✍ Brandt, Cuteri, Endrődi, Garnacho, Markó '22
- ▶ magnetic susceptibility
✍ Buividovich, Smith, von Smekal '21

beyond homogeneous magnetic fields

- ▶ inhomogeneous magnetic fields
✍ Valois et al. '21
- ▶ electric background fields
✍ Endrődi, Markó '22

Summary

Summary

- ▶ closing down on the $\mu_B > 0$ critical endpoint: lattice, FRG, DSE
- ▶ $T - \mu_I$ phase diagram (supersonic) pion condensation possible impact on cosmology
- ▶ $T - B$ phase diagram and the critical point

