

QCD matter with isospin-asymmetry

Gergely Endrődi

Goethe University of Frankfurt

in collaboration with Bastian Brandt and Sebastian Schmalzbauer



KHuK Tagung, Bad Honnef
1. December 2017

Outline

- ▶ introduction: QCD with isospin asymmetry
- ▶ setup
 - ▶ pion condensation
- ▶ results
 - ▶ phase boundary for pion condensation
 - ▶ deconfinement
 - ▶ QCD phase diagram
 - ▶ further applications
- ▶ conclusions

Introduction

Isospin chemical potential

- ▶ isospin density $n_I = n_u - n_d$
- ▶ $n_I < 0 \rightarrow$ excess of neutrons over protons
 \rightarrow excess of π^- over π^+



p^+



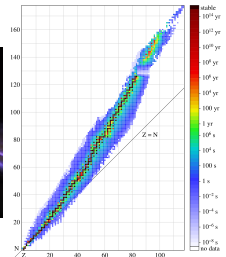
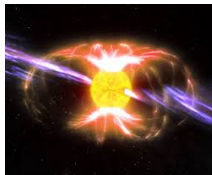
n



π^+

- ▶ applications

- ▶ neutron stars
- ▶ heavy-ion collisions (RHIC isobaric runs)



Isospin chemical potential

- ▶ isospin density $n_I = n_u - n_d$
- ▶ $n_I < 0 \rightarrow$ excess of neutrons over protons
 \rightarrow excess of π^- over π^+



p^+



n



π^+

- ▶ chemical potentials (2-flavor)

$$\mu_B = 3(\mu_u + \mu_d)/2 \qquad \mu_I = (\mu_u - \mu_d)/2$$

- ▶ here: zero baryon number but nonzero isospin

$$\mu_u = \mu_I \qquad \mu_d = -\mu_I$$

Methods

- ▶ QCD at low energies \approx pions
chiral perturbation theory

- ▶ on the level of charged pions: $\mu_\pi = 2\mu_I$

at zero temperature $\mu_\pi < m_\pi$

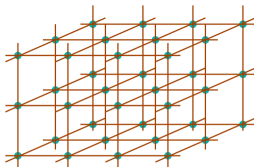
vacuum state

$$\mu_\pi = m_\pi$$

Bose-Einstein condensation

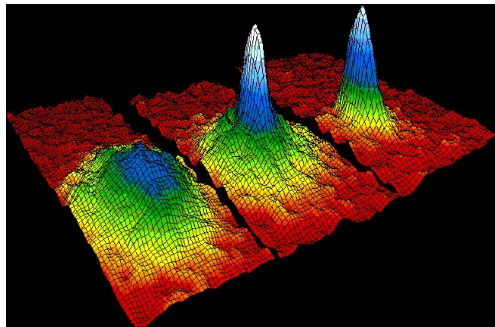
- ▶ on the level of quarks: lattice simulations

$$\mathcal{L}_{\text{QCD}} = \bar{\psi} M \psi + \frac{1}{4} \text{tr} F_{\mu\nu} F_{\mu\nu}$$



Bose-Einstein condensate

- ▶ accumulation of bosonic particles in lowest energy state

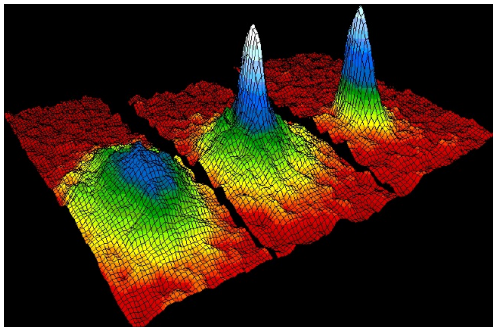


[Anderson et al '95 JILA-NIST/University of Colorado]

- ▶ velocity distribution of Ru atoms at low temperature
→ peak at zero velocity (zero energy)

Bose-Einstein condensate

- ▶ accumulation of bosonic particles in lowest energy state



[Anderson et al '95 JILA-NIST/University of Colorado]

- ▶ velocity distribution of Ru atoms at low temperature
→ peak at zero velocity (zero energy)
- ▶ phase transition, spontaneous symmetry breaking

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_l) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_l) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_I) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_I) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

- ▶ isospin chemical potential for the light quarks

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_I) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_I) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

- ▶ isospin chemical potential for the light quarks
- ▶ zero strangeness

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_I) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_I) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

- ▶ isospin chemical potential for the light quarks
- ▶ zero strangeness
- ▶ degenerate light quark masses

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_I) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_I) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

- ▶ isospin chemical potential for the light quarks
- ▶ zero strangeness
- ▶ degenerate light quark masses
- ▶ pionic source: explicit symmetry breaking

Lattice setup

- ▶ partition function

$$\mathcal{Z} = \int \mathcal{D}U \underbrace{\det \mathcal{M}_\ell}_{\text{light quarks}} \underbrace{\det \mathcal{M}_s}_{\text{strange quark}} \underbrace{e^{-S_g}}_{\text{gluons}}$$

$$\mathcal{M}_\ell = \begin{pmatrix} \not{D}(\mu_I) + m_\ell & \lambda \gamma_5 \\ -\lambda \gamma_5 & \not{D}(-\mu_I) + m_\ell \end{pmatrix} \quad \mathcal{M}_s = \not{D}(0) + m_s$$

- ▶ isospin chemical potential for the light quarks
- ▶ zero strangeness
- ▶ degenerate light quark masses
- ▶ pionic source: explicit symmetry breaking
necessary for spontaneous symmetry breaking in finite volume
needs to be extrapolated $\lambda \rightarrow 0$

Pion condensate from the lattice

Singular value representation

- ▶ pion condensate

$$\langle \pi^\pm \rangle = \frac{T}{V} \frac{\partial \log \mathcal{Z}}{\partial \lambda} = \dots = \frac{T}{V} \left\langle \text{Tr} \frac{2\lambda}{|\not{D}(\mu_l) + m_l|^2 + \lambda^2} \right\rangle$$

- ▶ singular values

$$|\not{D}(\mu_l) + m_l|^2 \psi_i = \xi_i^2 \psi_i$$

- ▶ spectral representation [Brandt, Endrődi 1611.06758]

$$\langle \pi^\pm \rangle = \frac{T}{V} \left\langle \sum_i \frac{2\lambda}{\xi_i^2 + \lambda^2} \right\rangle \xrightarrow{V \rightarrow \infty} \int d\xi \langle \rho(\xi) \rangle \frac{2\lambda}{\xi^2 + \lambda^2} \xrightarrow{\lambda \rightarrow 0} \pi \langle \rho(0) \rangle$$

first derived in [Kanazawa, Wettig, Yamamoto '11]

Singular value representation

- ▶ pion condensate

$$\langle \pi^\pm \rangle = \frac{T}{V} \frac{\partial \log \mathcal{Z}}{\partial \lambda} = \dots = \frac{T}{V} \left\langle \text{Tr} \frac{2\lambda}{|\not{D}(\mu_l) + m_l|^2 + \lambda^2} \right\rangle$$

- ▶ singular values

$$|\not{D}(\mu_l) + m_l|^2 \psi_i = \xi_i^2 \psi_i$$

- ▶ spectral representation [Brandt, Endrődi 1611.06758]

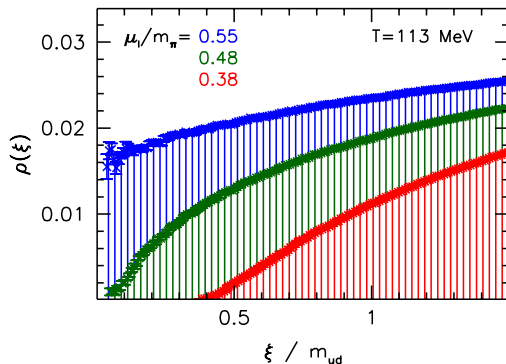
$$\langle \pi^\pm \rangle = \frac{T}{V} \left\langle \sum_i \frac{2\lambda}{\xi_i^2 + \lambda^2} \right\rangle \xrightarrow{V \rightarrow \infty} \int d\xi \langle \rho(\xi) \rangle \frac{2\lambda}{\xi^2 + \lambda^2} \xrightarrow{\lambda \rightarrow 0} \pi \langle \rho(0) \rangle$$

first derived in [Kanazawa, Wettig, Yamamoto '11]

- ▶ compare to Banks-Casher-relation at $\mu_l = 0$

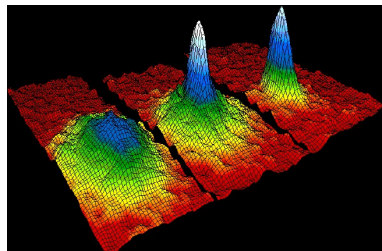
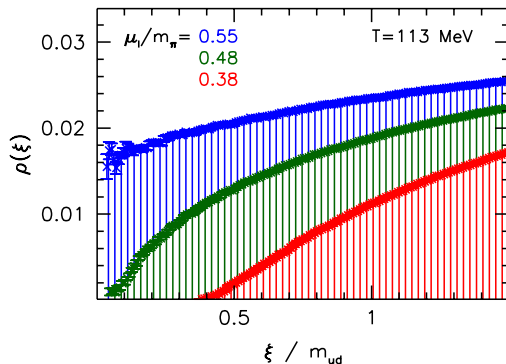
Singular value density

- ▶ spectral densities at $\lambda/m = 0.17$



Singular value density

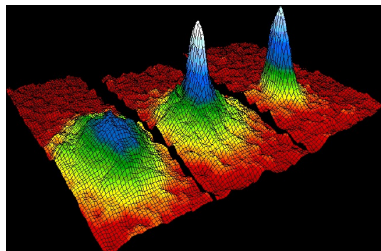
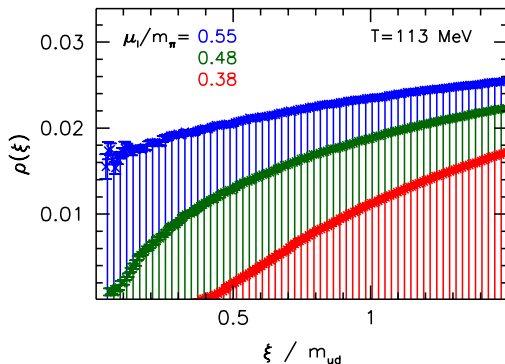
- ▶ spectral densities at $\lambda/m = 0.17$



- ▶ compare density of states around zero 'energy' ($\xi \approx 0$) to velocity distribution around zero

Singular value density

- ▶ spectral densities at $\lambda/m = 0.17$

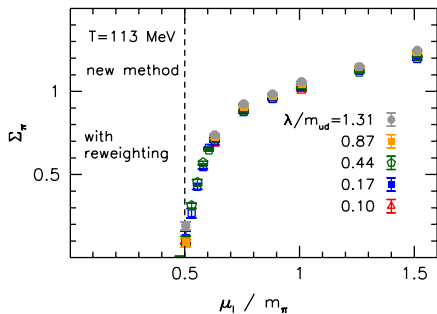


- ▶ compare density of states around zero 'energy' ($\xi \approx 0$) to velocity distribution around zero
- ▶ Bose-Einstein condensation!

Phase diagram

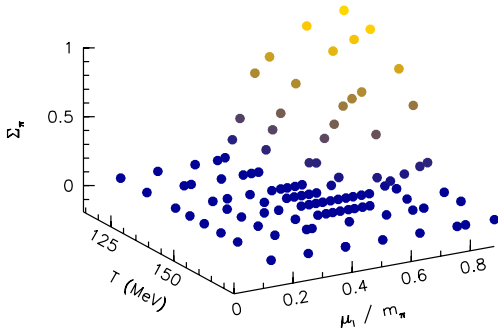
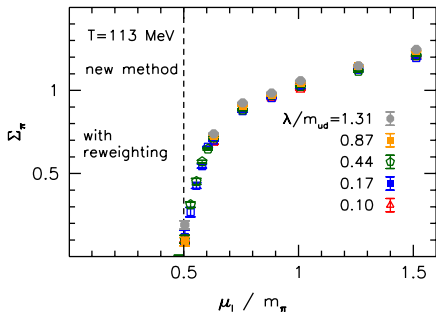
Pion condensate

- ▶ repeat this analysis for many different T and μ_I



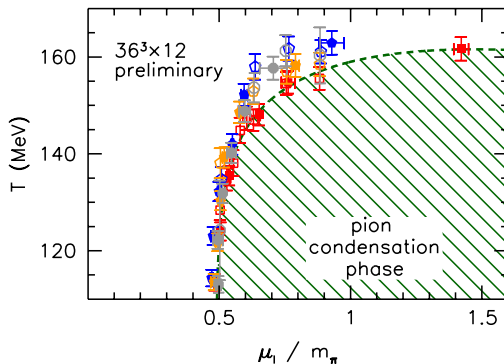
Pion condensate

- ▶ repeat this analysis for many different T and μ_I



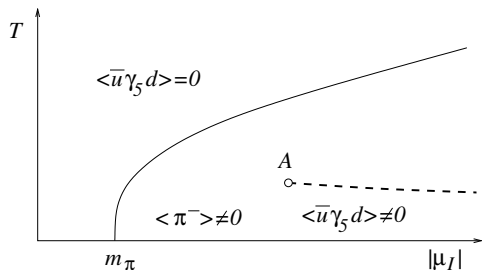
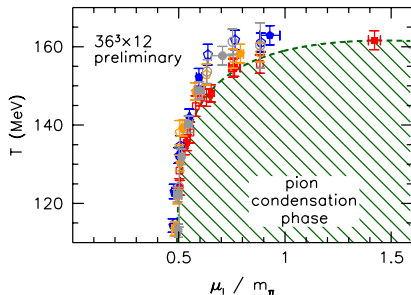
Phase boundary

- ▶ interpolate $\rho(0)$ as function of μ_I to find phase boundary



Phase boundary

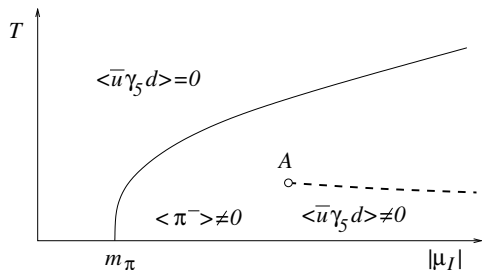
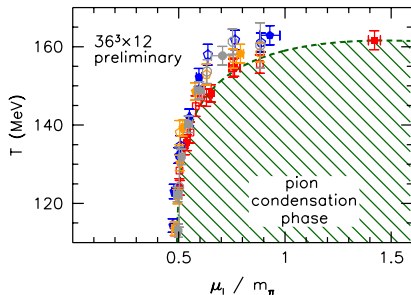
- ▶ interpolate $\rho(0)$ as function of μ_I to find phase boundary



- ▶ compare to expectations from χ PT [Son, Stephanov '00]

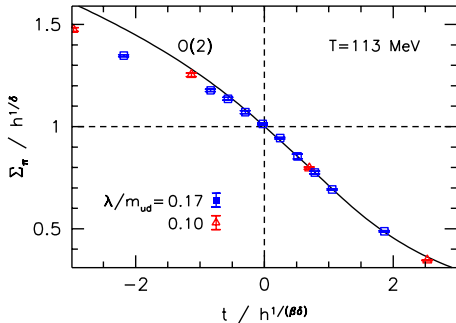
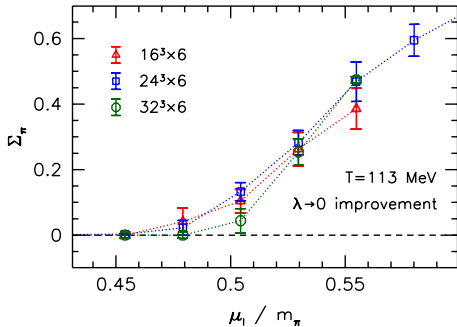
Phase boundary

- ▶ interpolate $\rho(0)$ as function of μ_I to find phase boundary



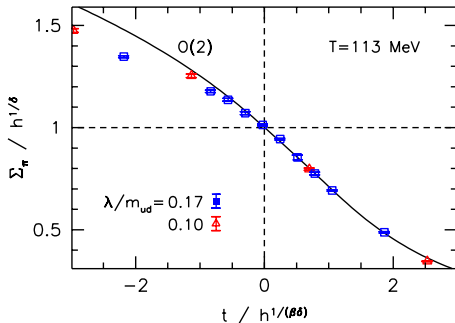
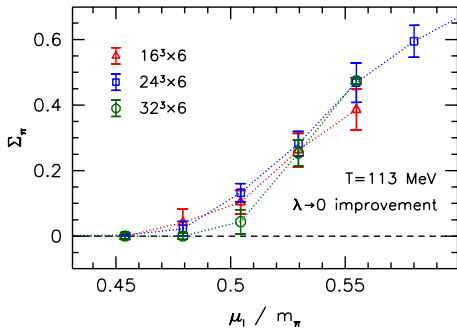
- ▶ compare to expectations from χ PT [Son, Stephanov '00]
- ▶ no pion condensate above $T \approx 160$ MeV [Brandt, Endrődi, Schmalzbauer 1709.10487]

Order of the transition - volume scaling



- ▶ volume scaling of order parameter shows typical sharpening
- ▶ collapse according to O(2) critical exponents [Ejiri et al '09]

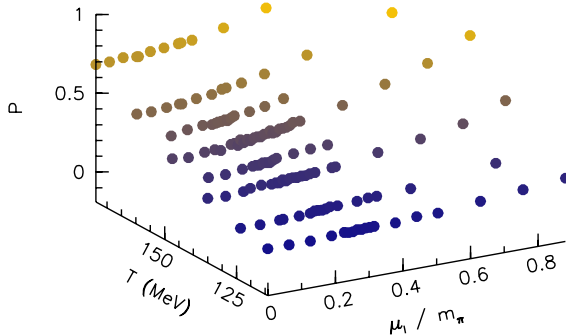
Order of the transition - volume scaling



- ▶ volume scaling of order parameter shows typical sharpening
- ▶ collapse according to $O(2)$ critical exponents [Ejiri et al '09]
- ▶ indications for a second order phase transition at $\mu_I = m_\pi/2$, in the $O(2)$ universality class

Deconfinement at high temperature

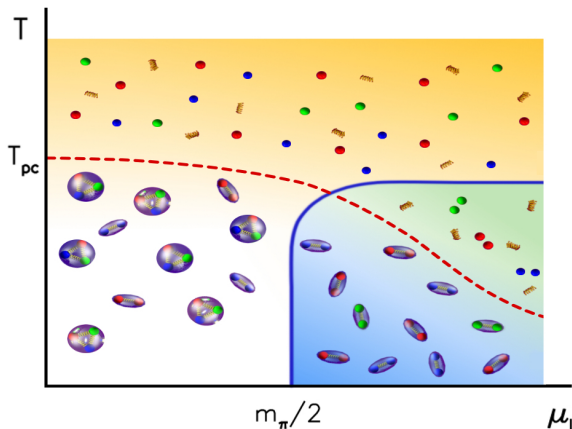
- ▶ deconfinement transition encoded in the Polyakov loop
 $P \sim \exp(-F_{q\bar{q}}/T)$



- ▶ deconfined matter for high μ_I : BCS superconductor with $u\bar{d}$ Cooper pairs [Son, Stephanov '02]

Phase diagram

- ▶ favored phase diagram schematically:
hadronic, quark-gluon plasma, BEC, BCS phases



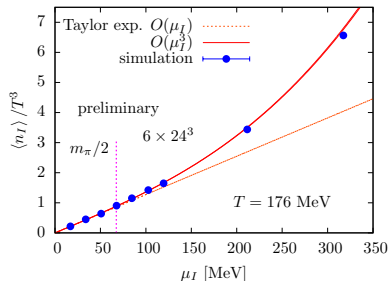
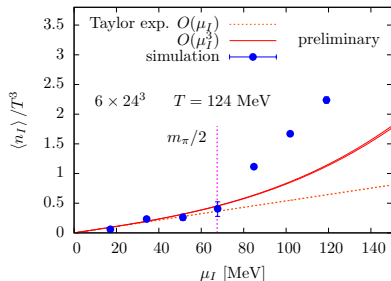
Further applications

Check Taylor-expansion

- ▶ isospin density via Taylor-expansion at $\mu_I = 0$

$$n_I(\mu_I) = \chi_2^I \cdot \mu_I + \chi_4^I \cdot \mu_I^3 + \dots$$

using $\chi_{2,4}^I$ from [BMWc, 1112.4416]



- ▶ low T : breakdown of expansion at $\mu_I = m_\pi/2$
- ▶ high T : pin down validity range of LO and NLO expansion

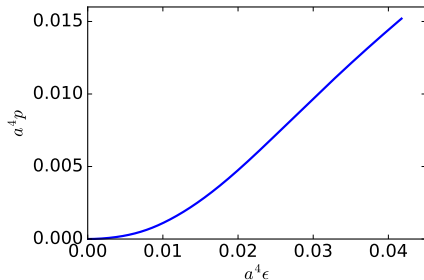
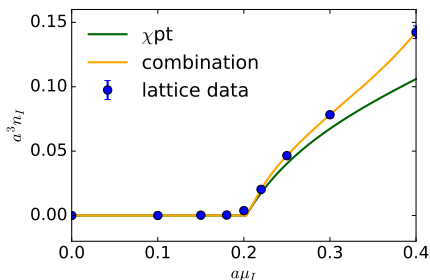
Equation of state

- ▶ pressure

$$p = \int_{m_\pi/2}^{\mu_l} d\mu'_l n_l(\mu'_l)$$

- ▶ energy density

$$\epsilon = -p + \mu_l n_l$$



- ▶ application: mass-radius relation of compact stars

Summary

- ▶ Bose-Einstein condensation via singular value density
 \rightsquigarrow flat extrapolation in λ
- ▶ established second-order phase transition at $\mu_I = m_\pi/2$
- ▶ QCD phase diagram with isospin asymmetry

