

# Causality violations in simulations of heavy ion collisions

Renata Krupczak – PhD student

University Bielefeld

Email: [rkrupczak@physik.uni-bielefeld.de](mailto:rkrupczak@physik.uni-bielefeld.de)



UNIVERSITÄT  
BIELEFELD



**Goal for this presentation:** Overview of causality conditions in hydrodynamics theory and understand the application of these conditions in heavy-ion collisions.

PHYSICAL REVIEW C **105**, L061901 (2022)

Letter

## Causality violations in realistic simulations of heavy-ion collisions

Christopher Plumberg <sup>1</sup>, Dekrayat Almaalol,<sup>1,2</sup> Travis Dore,<sup>1</sup> Jorge Noronha <sup>1</sup> and Jacquelyn Noronha-Hostler <sup>1</sup>

<sup>1</sup>*Illinois Center for Advanced Studies of the Universe, Department of Physics,  
University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA*

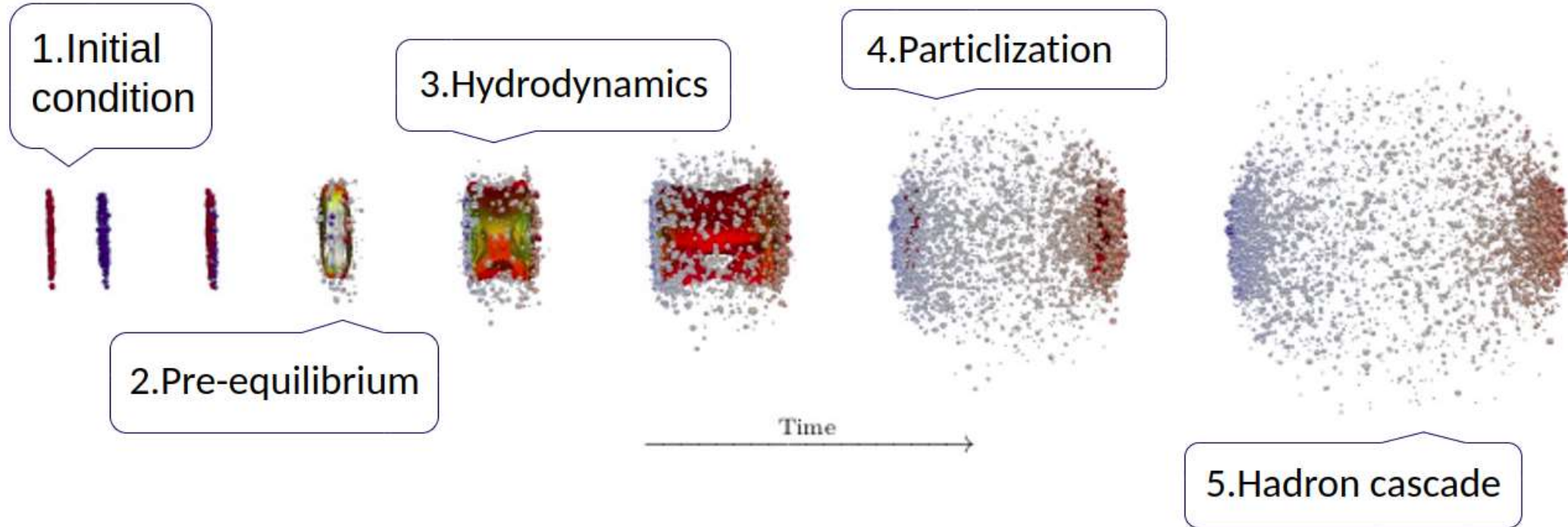
<sup>2</sup>*Department of Physics, Kent State University, Kent, Ohio 44242, USA*



(Received 3 May 2021; revised 4 March 2022; accepted 1 June 2022; published 22 June 2022)

Causality is violated in the early stages of state-of-the-art heavy-ion hydrodynamic simulations. Such violations are present in up to 75% of the fluid cells in the initial time and only after 2–3 fm/ $c$  of evolution do we find that 50% of the fluid cells are definitely causal. Superluminal propagation reaches up to 15% the speed of light in some of the fluid cells. The inclusion of pre-equilibrium evolution significantly reduces the number of acausal cells. Our findings suggests that relativistic causality may place constraints on the available parameter space of heavy-ion collision simulations when factored into more thorough statistical analyses.

# Simulation of heavy ion collision



Hydrodynamics to describe the QGP: relativistic, viscous and out of equilibrium

# Ideal Hydrodynamics

Perturbations

$$\epsilon = \epsilon_0 + \delta\epsilon$$

$$u^\mu = u_0^\mu + \delta u^\mu$$

$$u_\nu \nabla_\mu T_{(0)}^{\mu\nu} = 0$$

$$\Delta_\nu^\rho \nabla_\mu T_{(0)}^{\mu\nu} = 0$$

$$u_\nu \nabla_\mu \epsilon + (\epsilon + P) \nabla_\mu u^\mu = 0$$

$$(\epsilon + P) u^\mu \nabla_\mu u^\rho + c_s^2 \nabla_\mu \epsilon = 0$$

Use group velocity concept

$$v_g \equiv \left| \frac{d\omega}{dk} \right| = c_s \leq 1$$

Condition for causality

# Viscous Hydrodynamics – Navier Stokes (First order theory)

$$\nabla_{\mu} T^{\mu\nu} = 0$$

Shear and bulk viscosities

$$\begin{aligned} \pi^{\mu\nu} &= -\eta\sigma^{\mu\nu} \\ \Pi &= \zeta\lambda_{\lambda}^{\perp} u^{\lambda} \end{aligned}$$

$$v_g = \left| \frac{d\omega}{dk} \right| = 2\gamma_{\eta} k \text{ Grows without bound}$$

Unlimited velocity, acausal theory

Maxwell-Cataneo proposal

$$\tau_{\pi} D\pi^{\mu\nu} + \pi^{\mu\nu} = -\eta\sigma^{\mu\nu} + \dots$$

$$\tau_{\Pi} D\Pi + \Pi = \zeta\lambda_{\alpha}^{\perp} u^{\alpha} + \dots$$

Relaxation term

# Theories like Israel-Stewart: DNMR (Second order theory)

$$\tau_\pi \dot{\pi}^{\langle\mu\nu\rangle} + \pi^{\mu\nu} = 2\eta\sigma^{\mu\nu} + 2\tau_\pi \pi_\alpha^{\langle\mu} \omega^{\nu\rangle\alpha} - \delta_{\pi\pi} \pi^{\mu\nu} \theta + \phi_7 \pi_\alpha^{\langle\mu} \pi^{\nu\rangle\alpha} - \tau_{\pi\pi} \pi_\alpha^{\langle\mu} \sigma^{\nu\rangle\alpha} + \lambda_{\pi\Pi} \Pi \sigma^{\mu\nu} + \phi_6 \Pi \pi^{\mu\nu}$$

$$\tau_\Pi \dot{\Pi} + \Pi = -\zeta\theta - \delta_{\Pi\Pi} \Pi \theta + \phi_1 \Pi^2 + \lambda_{\Pi\pi} \pi^{\mu\nu} \sigma_{\mu\nu} + \phi_3 \pi^{\mu\nu} \pi_{\mu\nu}$$



Relaxation term

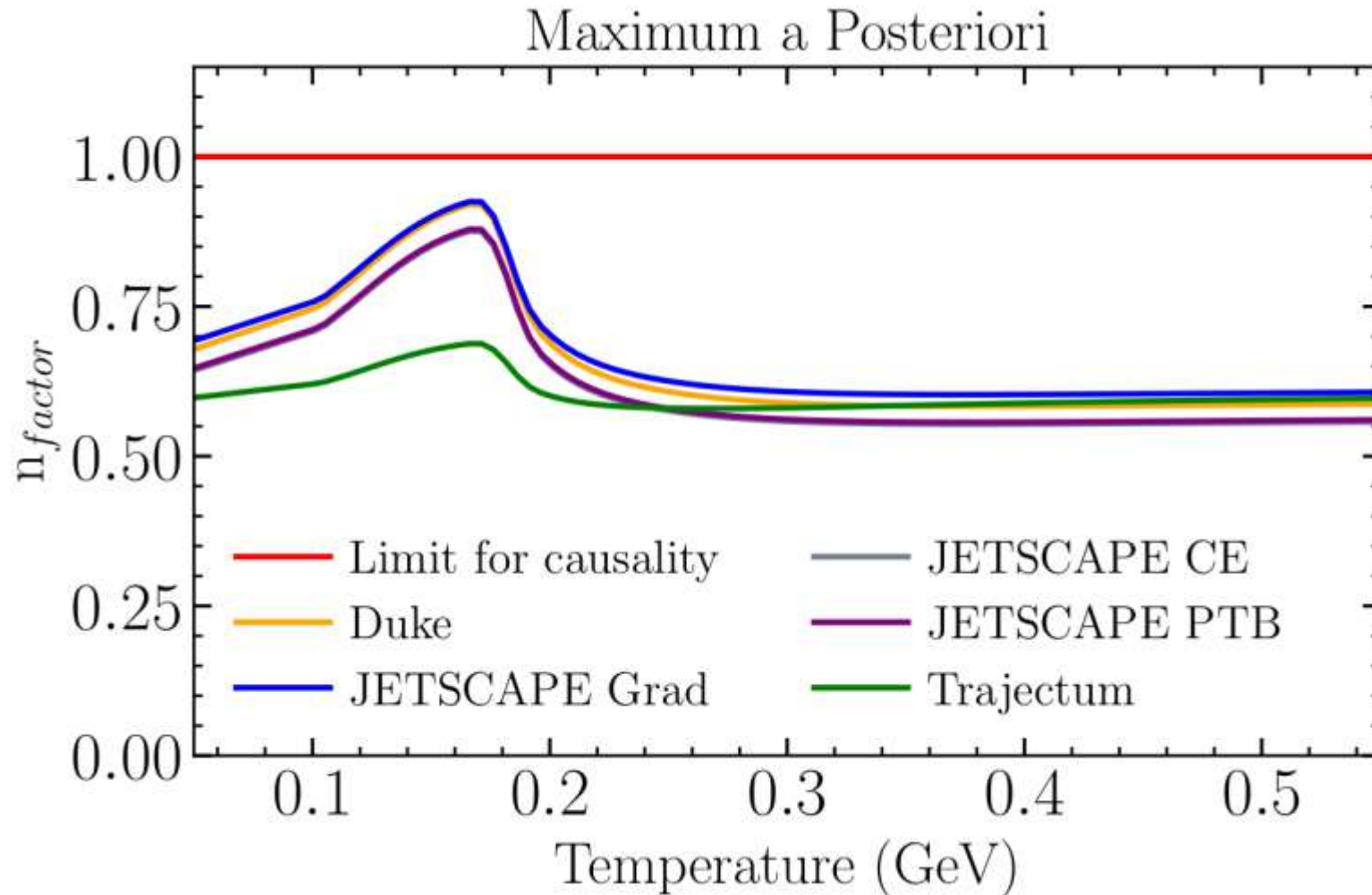
11 second order transport coefficients

$$\tau_\pi, \delta_{\pi\pi}, \phi_7, \tau_{\pi\pi}, \lambda_{\pi\Pi}, \phi_6, \tau_\Pi, \delta_{\Pi\Pi}, \phi_1, \lambda_{\Pi\pi}, \phi_3$$

Linear causality condition

$$c_s^2 + \frac{4}{3} \frac{\eta}{\tau_\pi (\epsilon + p)} + \frac{\zeta}{\tau_\Pi (\epsilon + p)} \leq 1$$

$$n_{\text{factor}} \equiv c_s^2 + \frac{4}{3} \frac{\eta}{\tau_\pi(\epsilon + p)} + \frac{\zeta}{\tau_\Pi(\epsilon + p)} \leq 1$$



# Theories like Israel-Stewart: DNMR

Nonlinear causality conditions are more complicated...

PHYSICAL REVIEW LETTERS **126**, 222301 (2021)

---

## Nonlinear Constraints on Relativistic Fluids Far from Equilibrium


Fábio S. Bemfica<sup>1,\*</sup>, Marcelo M. Disconzi<sup>2,†</sup>, Vu Hoang<sup>3,‡</sup>, Jorge Noronha<sup>4,§</sup> and Maria Radosz<sup>3,||</sup>

<sup>1</sup>*Escola de Ciências e Tecnologia, Universidade Federal do Rio Grande do Norte, 59072-970, Natal, Rio Grande do Norte, Brazil*

<sup>2</sup>*Department of Mathematics, Vanderbilt University, Nashville, Tennessee 37211, USA*

<sup>3</sup>*Department of Mathematics, The University of Texas at San Antonio, One UTSA Circle, San Antonio, Texas 78249, USA*

<sup>4</sup>*Illinois Center for Advanced Studies of the Universe, Department of Physics,  
University of Illinois at Urbana-Champaign, Urbana, Illinois 61801, USA*

 (Received 6 June 2020; revised 6 March 2021; accepted 29 April 2021; published 1 June 2021)

New constraints are found that must necessarily hold for Israel-Stewart-like theories of fluid dynamics to be causal far away from equilibrium. Conditions that are sufficient to ensure causality, local existence, and uniqueness of solutions in these theories are also presented. Our results hold in the full nonlinear regime, taking into account bulk and shear viscosities (at zero chemical potential), without any simplifying symmetry or near-equilibrium assumptions. Our findings provide fundamental constraints on the magnitude of viscous corrections in fluid dynamics far from equilibrium.

DOI: [10.1103/PhysRevLett.126.222301](https://doi.org/10.1103/PhysRevLett.126.222301)



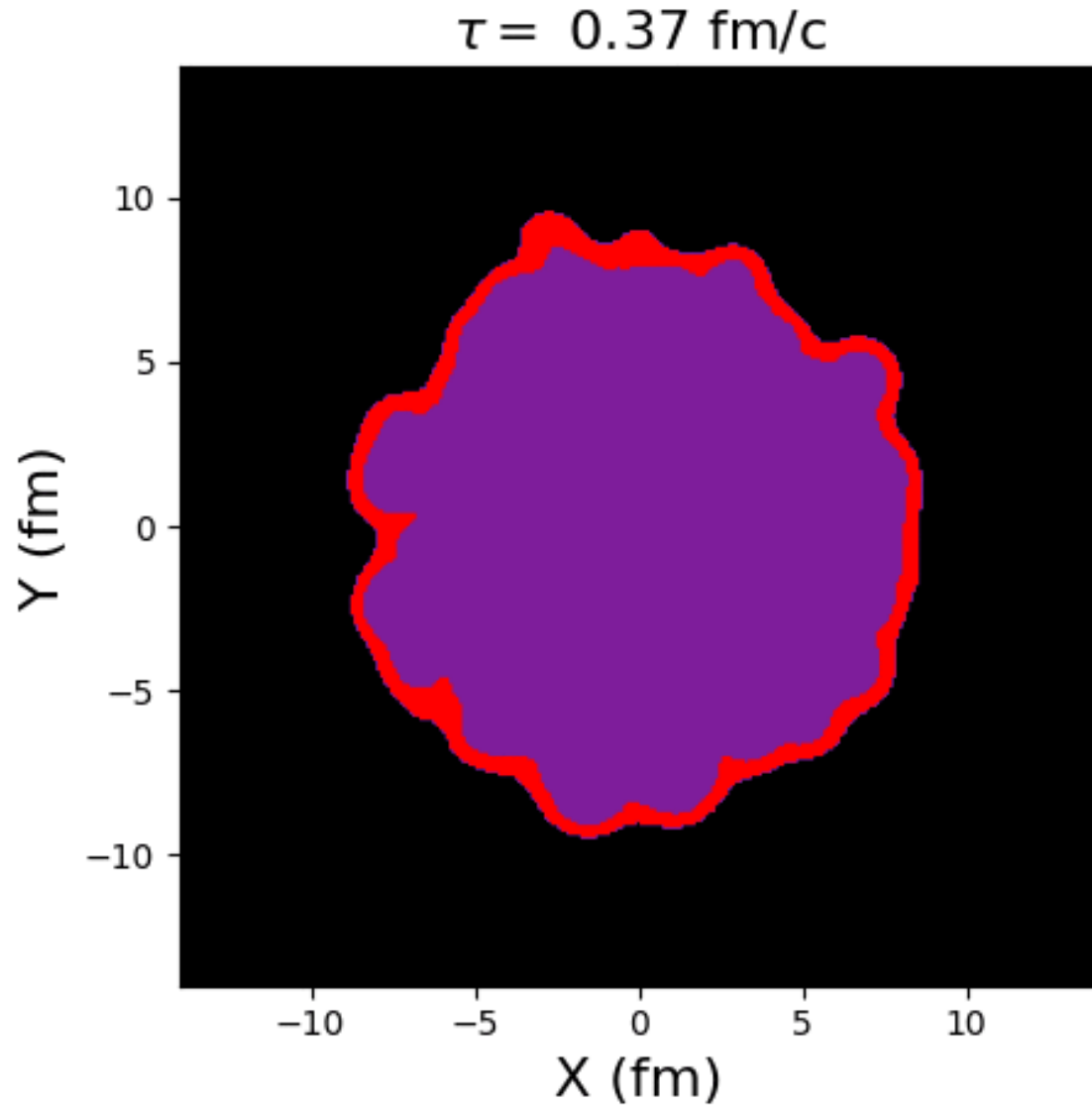
## Necessary conditions

$$\begin{aligned}
 & (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{1}{2}\tau_{\pi\pi} |\Lambda_1| \geq 0, \\
 & \varepsilon + P + \Pi - \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{4\tau_\pi} \Lambda_3 \geq 0, \\
 & \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) + \frac{\tau_{\pi\pi}}{4\tau_\pi} (\Lambda_a + \Lambda_d) \geq 0, \quad a \neq d, \\
 & \varepsilon + P + \Pi + \Lambda_a - \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{4\tau_\pi} (\Lambda_d + \Lambda_a) \geq 0, \quad a \neq d \\
 & \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) + \frac{\tau_{\pi\pi}}{2\tau_\pi} \Lambda_d + \frac{1}{6\tau_\pi} [2\eta + \lambda_{\pi\Pi\Pi} + (6\delta_{\pi\pi} - \tau_{\pi\pi}) \Lambda_d] \\
 & \quad + \frac{\zeta + \delta_{\Pi\Pi\Pi} + \lambda_{\Pi\Pi} \Lambda_d}{\tau_\Pi} + (\varepsilon + P + \Pi + \Lambda_d) c_s^2 \geq 0 \\
 & \varepsilon + P + \Pi + \Lambda_d - \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{2\tau_\pi} \Lambda_d - \frac{1}{6\tau_\pi} [2\eta + \lambda_{\pi\Pi\Pi} + (6\delta_{\pi\pi} - \tau_{\pi\pi}) \Lambda_d] \\
 & \quad - \frac{\zeta + \delta_{\Pi\Pi\Pi} + \lambda_{\Pi\Pi} \Lambda_d}{\tau_\Pi} - (\varepsilon + P + \Pi + \Lambda_d) c_s^2 \geq 0
 \end{aligned}$$

## Sufficient conditions

$$\begin{aligned}
 & (\varepsilon + P + \Pi - |\Lambda_1|) - \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{2\tau_\pi} \Lambda_3 \geq 0 \\
 & (2\eta + \lambda_{\pi\Pi\Pi}) - \tau_{\pi\pi} |\Lambda_1| > 0, \\
 & \tau_{\pi\pi} \leq 6\delta_{\pi\pi}, \\
 & \frac{\lambda_{\Pi\Pi}}{\tau_\Pi} + c_s^2 - \frac{\tau_{\pi\pi}}{12\tau_\pi} \geq 0, \\
 & \frac{1}{3\tau_\pi} [4\eta + 2\lambda_{\pi\Pi\Pi} + (3\delta_{\pi\pi} + \tau_{\pi\pi}) \Lambda_3] + \frac{\zeta + \delta_{\Pi\Pi\Pi} + \lambda_{\Pi\Pi} \Lambda_3}{\tau_\Pi} + |\Lambda_1| + \Lambda_3 c_s^2 \\
 & \quad + \frac{\frac{12\delta_{\pi\pi} - \tau_{\pi\pi}}{12\tau_\pi} \left( \frac{\lambda_{\Pi\Pi}}{\tau_\Pi} + c_s^2 - \frac{\tau_{\pi\pi}}{12\tau_\pi} \right) (\Lambda_3 + |\Lambda_1|)^2}{\varepsilon + P + \Pi - |\Lambda_1| - \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{2\tau_\pi} \Lambda_3} \leq (\varepsilon + P + \Pi) (1 - c_s^2), \\
 & \frac{1}{6\tau_\pi} [2\eta + \lambda_{\pi\Pi\Pi} + (\tau_{\pi\pi} - 6\delta_{\pi\pi}) |\Lambda_1|] + \frac{\zeta + \delta_{\Pi\Pi\Pi} - \lambda_{\Pi\Pi} |\Lambda_1|}{\tau_\Pi} + (\varepsilon + P + \Pi - |\Lambda_1|) c_s^2 \geq 0, \\
 & 1 \geq \frac{\frac{12\delta_{\pi\pi} - \tau_{\pi\pi}}{12\tau_\pi} \left( \frac{\lambda_{\Pi\Pi}}{\tau_\Pi} + c_s^2 - \frac{\tau_{\pi\pi}}{12\tau_\pi} \right) (\Lambda_3 + |\Lambda_1|)^2}{\left[ \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) - \frac{\tau_{\pi\pi}}{2\tau_\pi} |\Lambda_1| \right]^2}, \\
 & \frac{1}{3\tau_\pi} [4\eta + 2\lambda_{\pi\Pi\Pi} - (3\delta_{\pi\pi} + \tau_{\pi\pi}) |\Lambda_1|] + \frac{\zeta + \delta_{\Pi\Pi\Pi} - \lambda_{\Pi\Pi} |\Lambda_1|}{\tau_\Pi} + (\varepsilon + P + \Pi - |\Lambda_1|) c_s^2 \\
 & \geq \frac{(\varepsilon + P + \Pi + \Lambda_2) (\varepsilon + P + \Pi + \Lambda_3)}{3(\varepsilon + P + \Pi - |\Lambda_1|)} \times \left\{ 1 + \frac{2 \left[ \frac{1}{2\tau_\pi} (2\eta + \lambda_{\pi\Pi\Pi}) + \frac{\tau_{\pi\pi}}{2\tau_\pi} \Lambda_3 \right]}{\varepsilon + P + \Pi - |\Lambda_1|} \right\},
 \end{aligned}$$

# Testing the causality conditions in the heavy ion simulations



## CAUSAL

Necessary ✓

Sufficient ✓

## INDETERMINATE

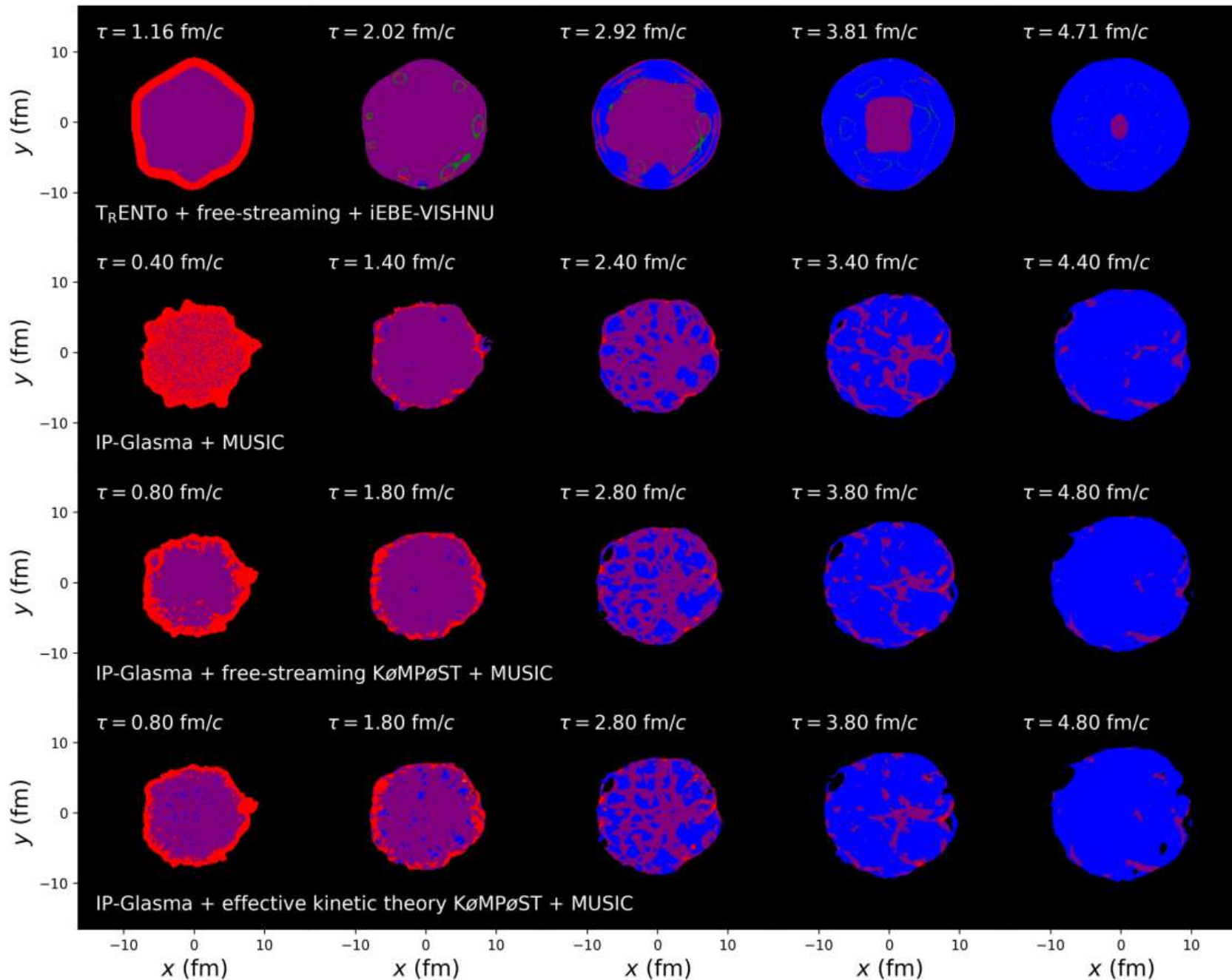
Necessary ✓

Sufficient ✗

## ACAUSAL

Sufficient ✗

Necessary ✗



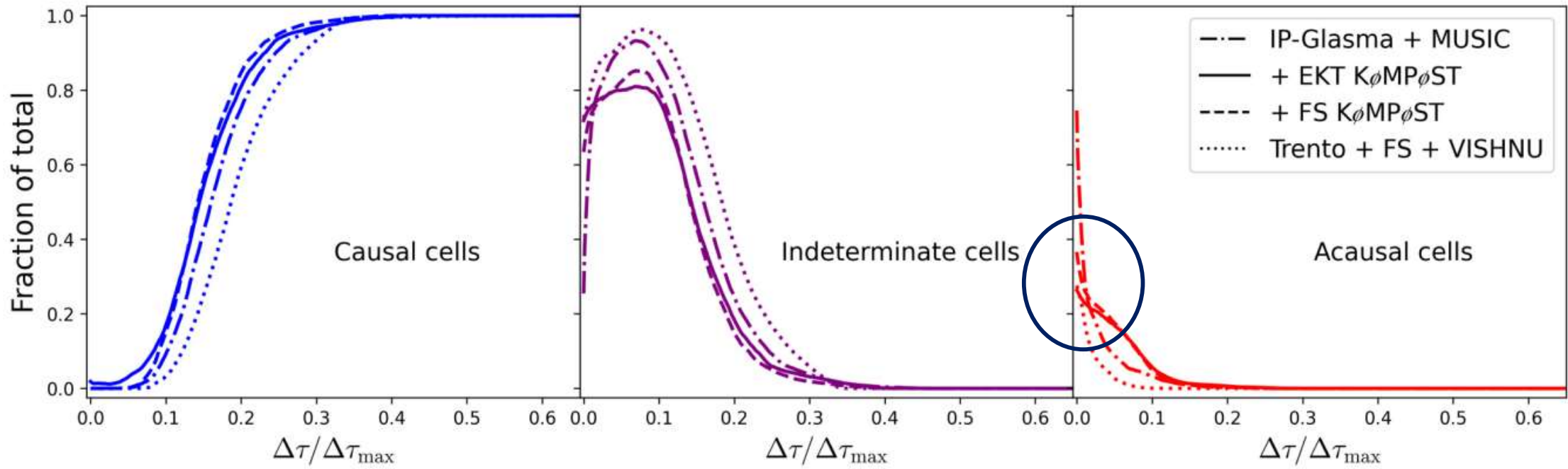
**ACAUSAL**

**INDETERMINATE**

**CAUSAL**

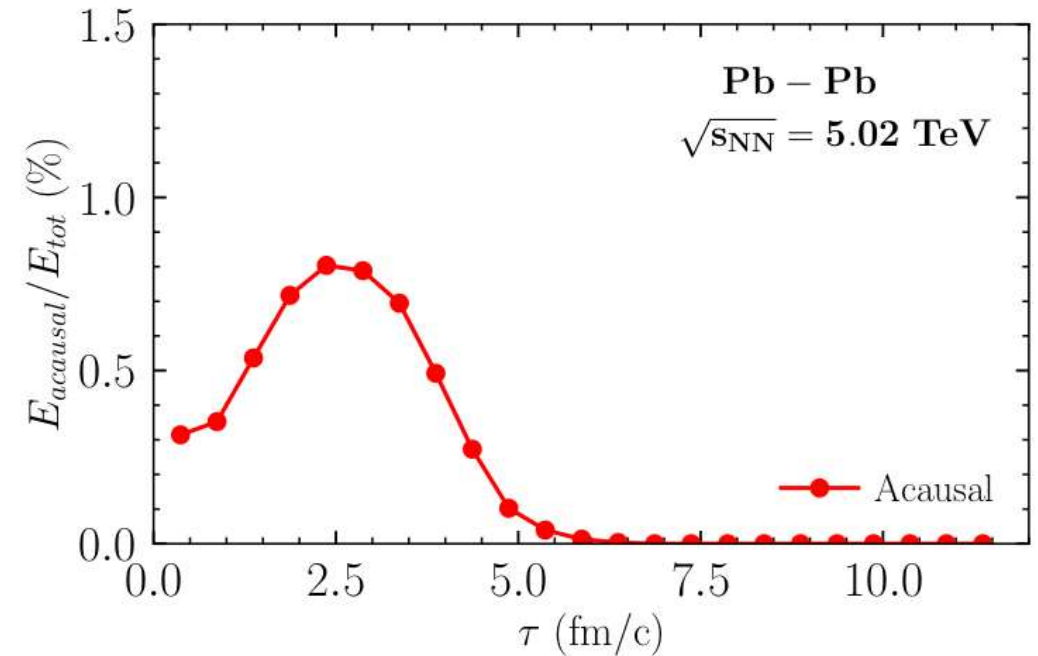
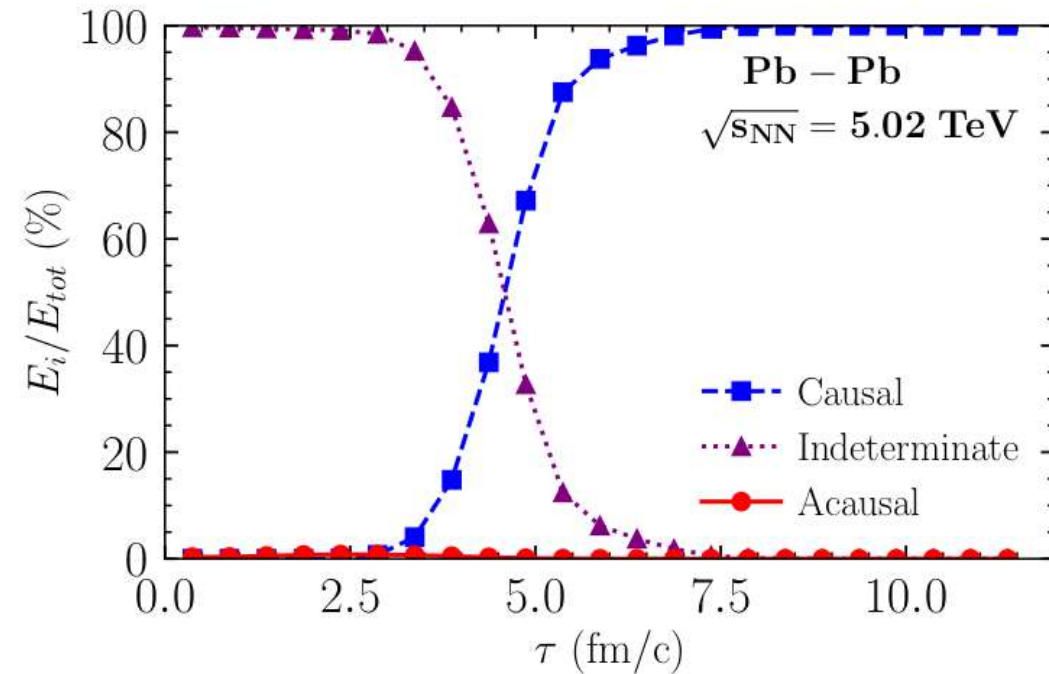
Conclusion:  
causality violation  
seems to be always  
present

PLUMBERG, Christopher et al. Causality violations in realistic simulations of heavy-ion collisions. **Physical Review C**, v. 105, n. 6, p. L061901, 2022.



Pre-equilibrium models can significantly reduce the fraction of acausal cells in the system.

In terms of energy, it represents a small part of the system.



Trento + Free streaming + MUSIC

Studying the conditions for causality, the most often and important condition:

$$\epsilon + P + \Pi + \Lambda_d - \frac{1}{2\tau_\pi}(2\eta + \lambda_{\pi\Pi}\Pi) - \frac{\tau_{\pi\pi}}{2\tau_\pi}\Lambda_d - \frac{1}{6\tau_\pi}[2\eta + \lambda_{\pi\Pi}\Pi + (6\delta_{\pi\pi} - \tau_{\pi\pi})\Lambda_d] - \frac{\zeta + \delta_{\Pi\Pi}\Pi + \lambda_{\Pi\pi}\Lambda_d}{\tau_\Pi} - (\epsilon + P + \Pi + \Lambda_d)c_s^2 \geq 0$$

Contains the linear causality condition inside:

$$c_s^2 + \frac{4}{3} \frac{\eta}{\tau_\pi(\epsilon + p)} + \frac{\zeta}{\tau_\Pi(\epsilon + p)} \leq 1$$

Studying the conditions for causality, the most often and important condition:

$$\epsilon + P + \Pi + \Lambda_d - \frac{1}{2\tau_\pi}(2\eta + \lambda_{\pi\Pi}\Pi) - \frac{\tau_{\pi\pi}}{2\tau_\pi}\Lambda_d - \frac{1}{6\tau_\pi}[2\eta + \lambda_{\pi\Pi}\Pi + (6\delta_{\pi\pi} - \tau_{\pi\pi})\Lambda_d]$$

$$\frac{-\zeta + \delta_{\Pi\Pi}\Pi + \lambda_{\Pi\pi}\Lambda_d}{\tau_{\Pi}} - (\epsilon + P + \Pi + \Lambda_d)c_s^2 \geq 0$$

When this term is equal to zero, the acausality disappears from the system.

$$\frac{-\zeta}{\tau_{\Pi}} \rightarrow 0 \quad \longrightarrow \quad \tau_{\Pi} = b_{\Pi} \frac{\zeta}{(1/3 - c_s^2)^2 (\epsilon + P)}$$

Increase the bulk relaxation time

Studying the conditions for causality, the most often and important condition:

$$\epsilon + P + \Pi + \Lambda_d - \frac{1}{2\tau_\pi}(2\eta + \lambda_{\pi\Pi}\Pi) - \frac{\tau_{\pi\pi}}{2\tau_\pi}\Lambda_d - \frac{1}{6\tau_\pi}[2\eta + \lambda_{\pi\Pi}\Pi + (6\delta_{\pi\pi} - \tau_{\pi\pi})\Lambda_d]$$

$$\frac{-\zeta + \delta_{\Pi\Pi}\Pi + \lambda_{\Pi\pi}\Lambda_d}{\tau_\Pi} - (\epsilon + P + \Pi + \Lambda_d)c_s^2 \geq 0$$

When this term is equal to zero, the acausality disappears from the system.

$$\frac{-\zeta}{\tau_\Pi} \rightarrow 0 \quad \longrightarrow \quad \tau_\Pi = b_\Pi \frac{\zeta}{(1/3 - c_s^2)^2 (\epsilon + P)}$$

Increase this parameter

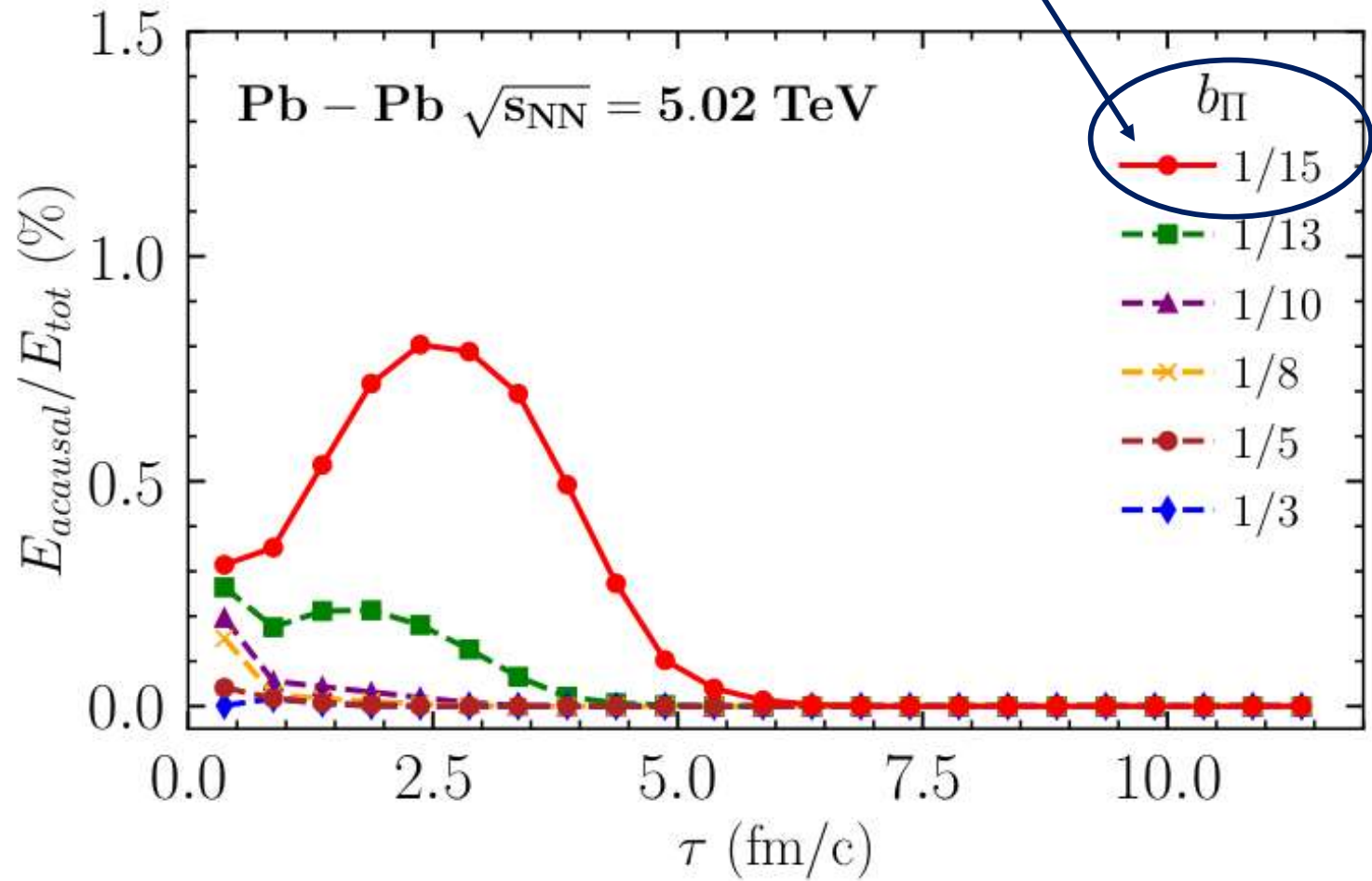
Properties:

- Dimensionless parameter
- Without physical meaning
- It is not fixed by hydrodynamic theory
- This change doesn't affect the final observables

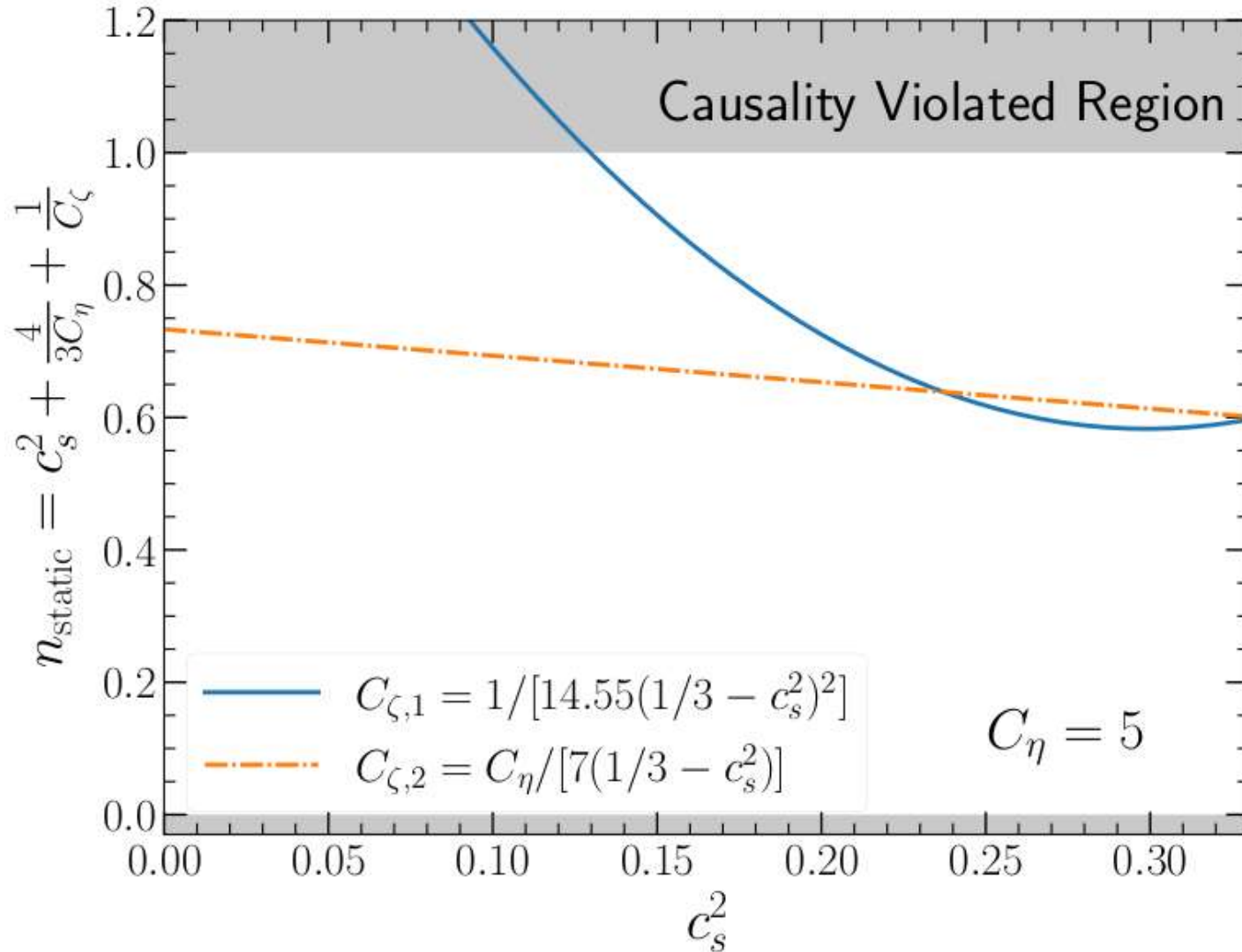


Standard value used in MUSIC (DNMR theory)

Increase this parameter, you can reduce significantly the acausality



# Parameters related with the causality condition



$$C_\zeta \propto \frac{1}{(1/3 - c_s^2)^\alpha}$$

$$C_{\zeta,1} = \frac{b_\Pi}{(1/3 - c_s^2)^2} \longrightarrow \begin{cases} b_\Pi = 1/14.55 \\ \alpha = 2 \end{cases}$$

Kinetic approach

$$C_{\zeta,2} = \frac{b_\Pi b_\pi}{(1/3 - c_s^2)} \longrightarrow \begin{cases} b_\Pi = 1/7 \\ \alpha = 1 \end{cases}$$

Strongly coupled theory

The sound speed that comes from the equation of state has a minimum value  $\sim 0.15$

## Conclusions...

- Acausality in heavy ion simulations is a recent topic with many open questions;
- Causality is violated in linear regime in some specific conditions, but can be controlled;
- In nonlinear regimes, the causality condition can't affirm if the system is only causal or acausal;
- Causality violations exist in heavy ion collisions independent of the computational model;
- Parametrizations for the hydrodynamics, as for example the relaxation time plays an important rule;
- The causality violation does not appear to affect significantly the final observables, which makes the issue more complex to solve.

Thank you!